

Optimal quantum circuit cuts

(with application to clustered Hamiltonian simulation)

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Overview

I. **Background & summary of results**

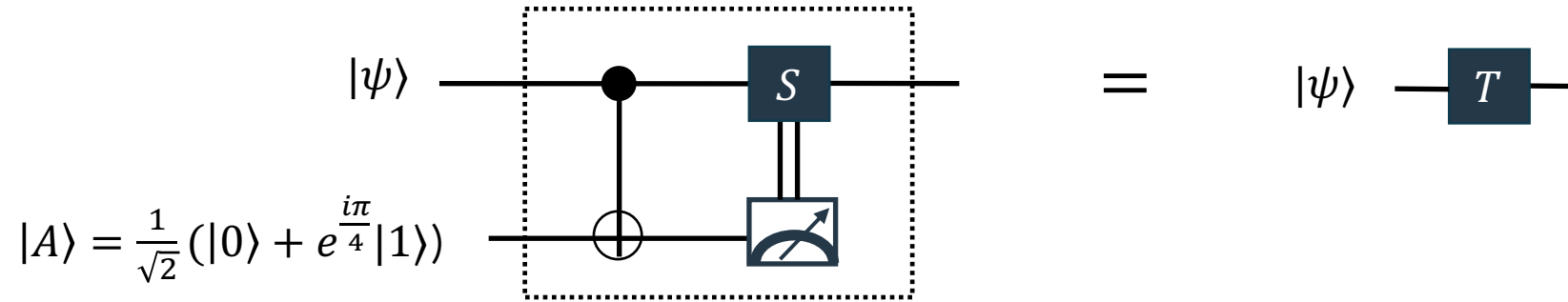
II. Optimal space-like cuts

III. Optimal time-like cuts

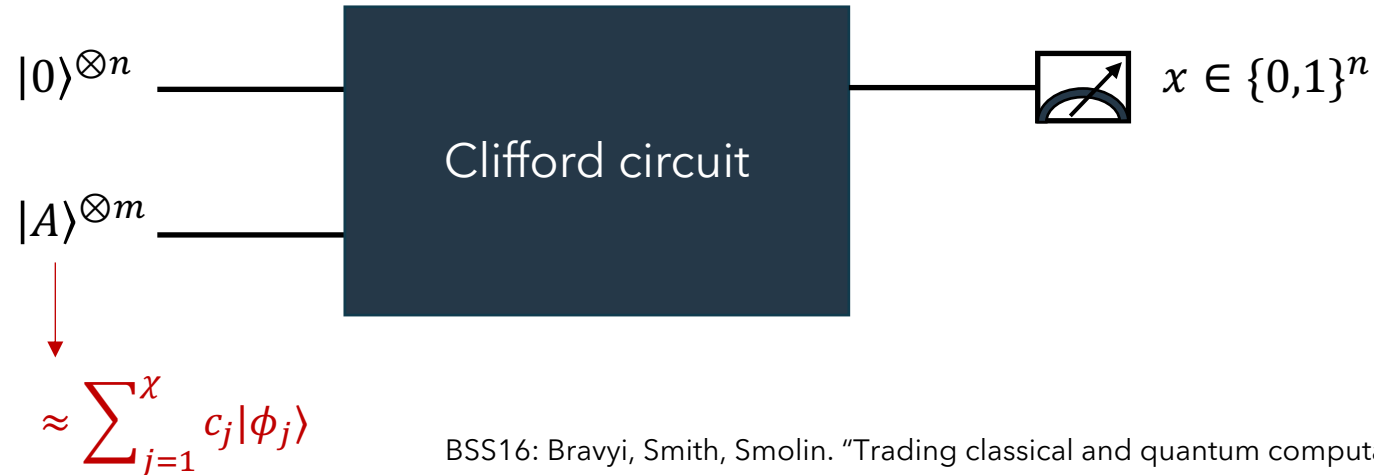
IV. Lower bounds

Quantum information processing consumes quantum **resources**. E.g., **entanglement** & **magic**.

Fault-tolerant quantum computation with **magic** states:



Replace **magic** with random Clifford operations [BSS16, Bra+19]:

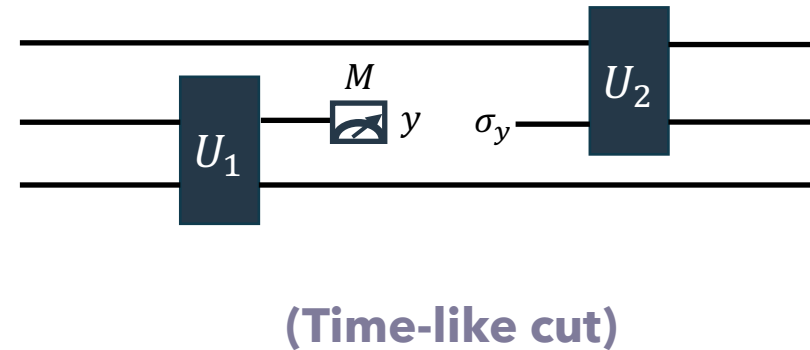
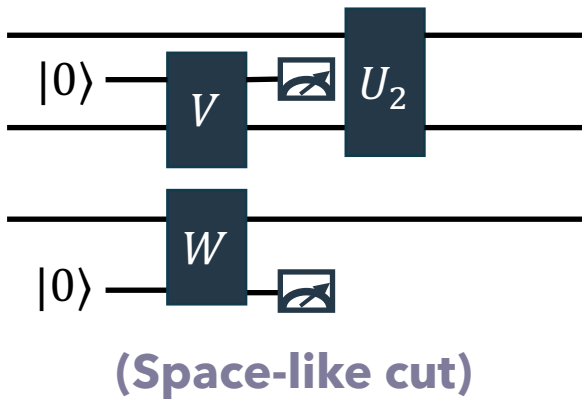
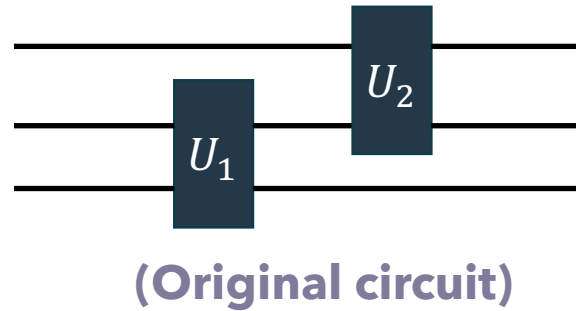


BSS16: Bravyi, Smith, Smolin. "Trading classical and quantum computational resources" PRX 6.2 (2016)

Bra+19: Bravyi et al. "Simulation of quantum circuits by low-rank stabilizer decompositions" Quantum 3 (2019)

Q1: Can we replace **entanglement** with random local operations?

→ If the goal is to compute expectation values, then yes.

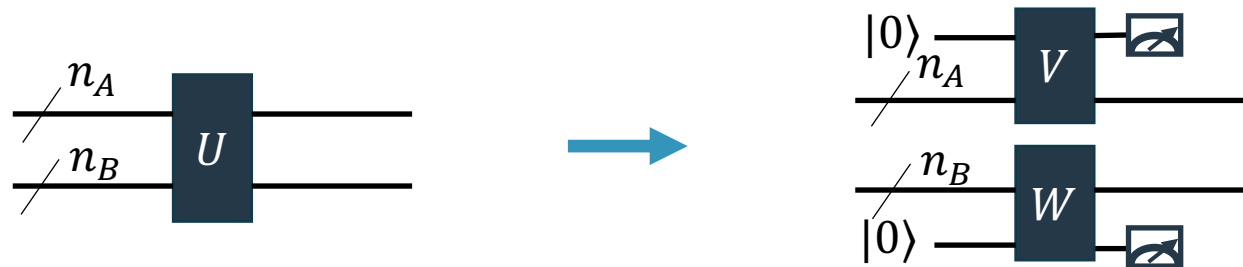


V , W , and M are potentially random.

Q2: At what cost?

Summary of results

For **space-like cuts** the cost is related to the **entangling power**.



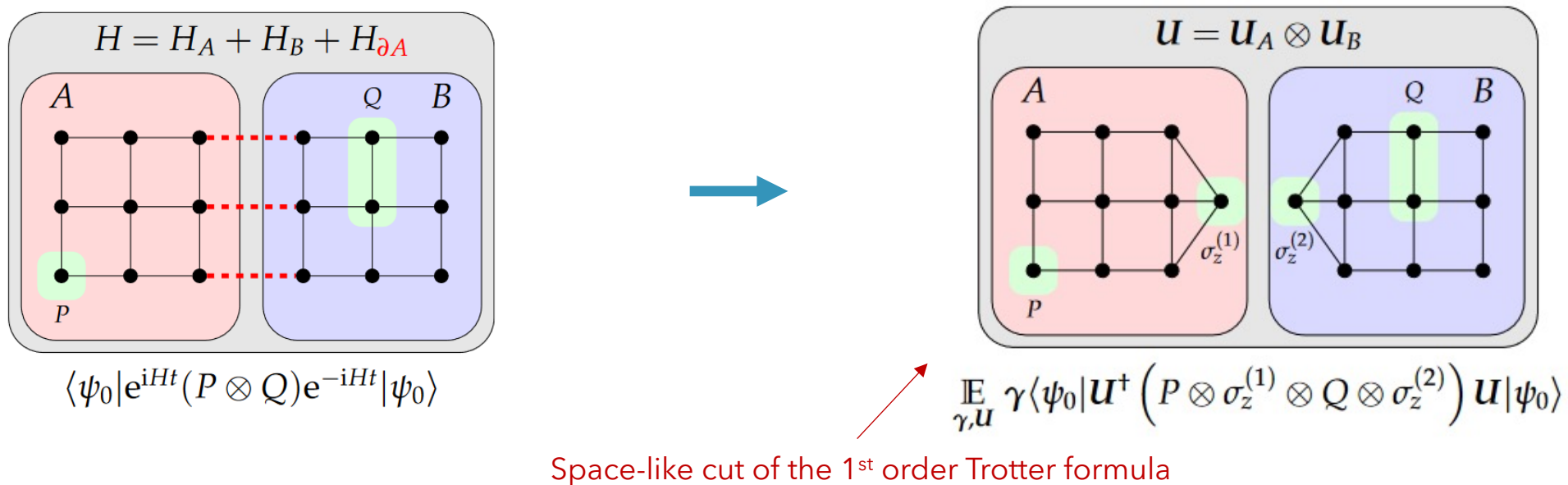
Entanglement measures

	CC?	Cost	Ancillas	Restrictions
** [BSS16]	No	$2^{O(n_B)}$	1	$O(1)$ -sparse
[PS22]	Yes	$1 + 2R(J(\mathcal{U}))$	$n_A + n_B$	Clifford
[This work]	No	$\xi(\mathcal{U})$	2	None

We have $1 + 2R(J(\mathcal{U})) \leq \xi(\mathcal{U}) \leq 2^{2(n_A+n_B)+1} - 1$. In many cases, $1 + 2R(J(\mathcal{U})) = \xi(\mathcal{U})$

Summary of results

Space-like cuts can be applied to **remove interactions** in a Hamiltonian simulation.

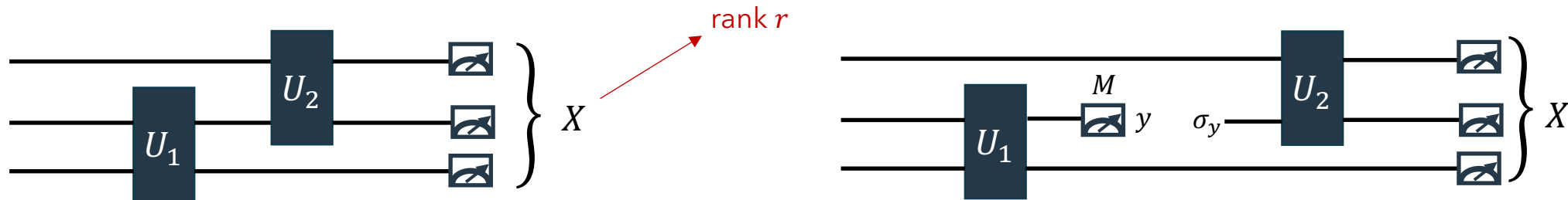


Runtime is on the order of $\exp(8\eta t)$ where $\eta = \sum_{k \in \partial A} \|H_k\|$.

Best prior work achieved $\exp(|\partial A| \eta^{1/p} t^{1+1/p} / \epsilon^{1/p})$ using a p^{th} order formula [Chi+21].

Summary of results

For **time-like cuts** the cost depends on the **number of wires** replaced.



	CC?	Cost	Ancillas
[Pen+19]	No	16^k	None
[BPS23]	Yes	4^k	$2k$
[This work]	Yes	$2^k r$	None

Lower bound also obtained: $\Theta(2^k)$ samples are necessary and sufficient to estimate output probabilities.

Pen+19: Peng et al. "Simulating Large Quantum Circuits on a Small Quantum Computer". PRL 125 (2020)

BPS23: Brenner, Piveteau, and Sutter. Optimal wire cutting with classical communication. arxiv:1506.01396

Overview

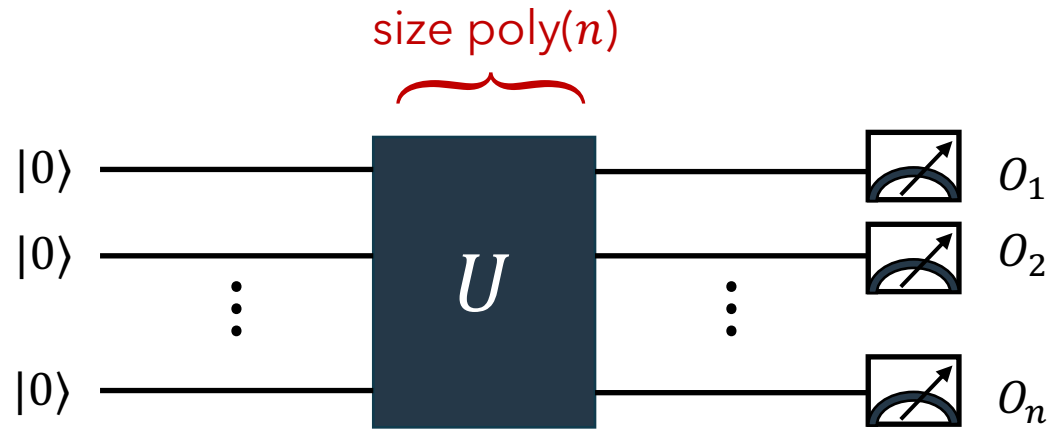
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Computational model



Goal: Compute a good estimate of $\langle O_1 \otimes O_2 \otimes \cdots \otimes O_n \rangle = \langle 0 | U^\dagger (O_1 \otimes O_2 \otimes \cdots \otimes O_n) U | 0 \rangle$.

Standard solution: Run the circuit M times and compute the empirical mean,

$$\mu = \frac{1}{M} (\mu_1 + \mu_2 + \cdots + \mu_M).$$

$$\text{Var}[\mu] \sim \frac{1}{M}$$

A red arrow points from the text "Standard solution:" to this equation.

Quasiprobability-based simulation

A **quasiprobability decomposition (QPD)** of a channel $\mathcal{N}_{A \rightarrow A}$ has the form

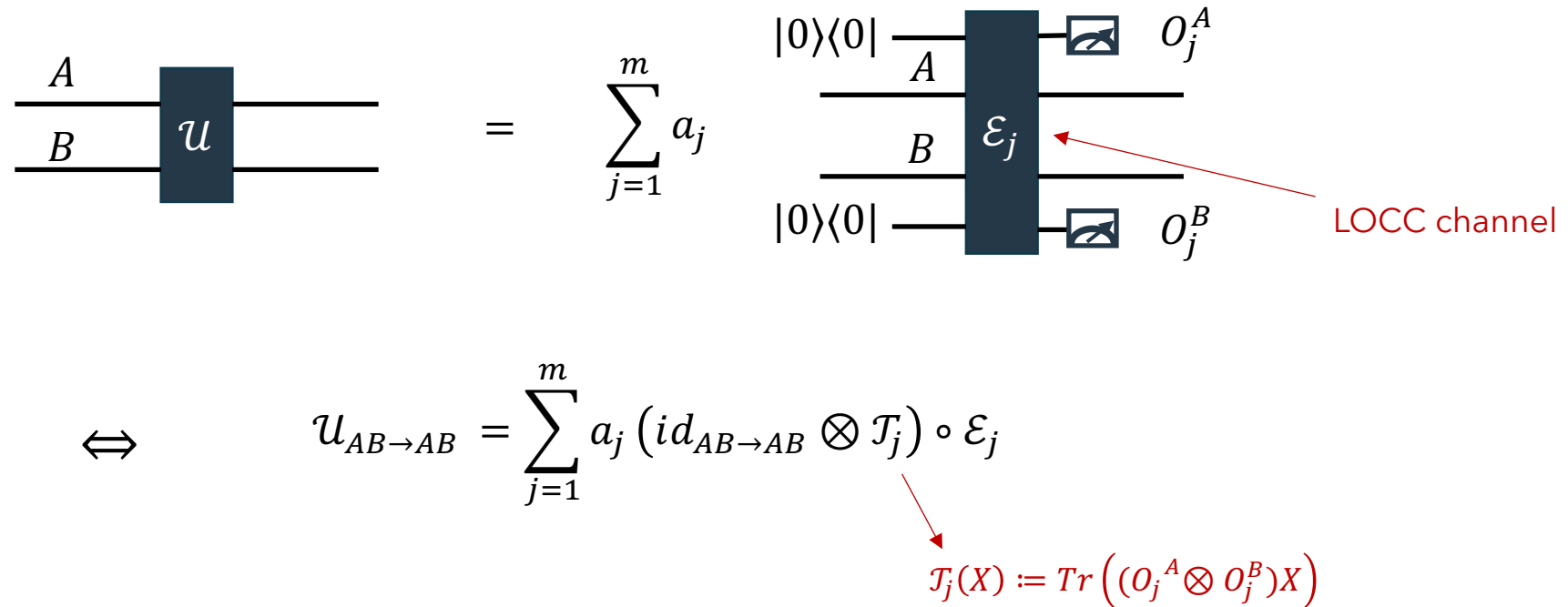
$$\mathcal{N} = \sum_{j=1}^m a_j \mathcal{T}_j \circ \mathcal{E}_j = \|a\|_1 \sum_{j=1}^m \frac{|a_j|}{\|a\|_1} \text{sign}(a_j) \mathcal{T}_j \circ \mathcal{E}_j$$

$$\Leftrightarrow \begin{array}{c} A \\ \text{---} \end{array} \boxed{\mathcal{N}} \begin{array}{c} A \\ \text{---} \end{array} = \sum_{j=1}^m a_j \begin{array}{c} A \\ \text{---} \end{array} \boxed{\mathcal{E}_j} \begin{array}{c} A \\ \text{---} \\ R \\ \text{---} \end{array} \boxed{\text{meter}} O_j$$

The 1-norm of the QPD is defined as $\|a\|_1 = \sum_{j=1}^m |a_j|$.

Fact: Let X be a Hermitian observable on A . Measuring the observable $\|a\|_1 \text{sign}(a_j) (O_j \otimes X)$ on the ensemble of states $\left\{ \left(|a_j| / \|a\|_1, \mathcal{E}_j(\rho) \right) \right\}_j$ yields an unbiased estimator of $\text{Tr}(X\mathcal{N}(\rho))$ for any ρ .

A **space-like cut** is a QPD of the form



Theorem [PS22]: The minimal 1-norm of a space-like cut of $\mathcal{U}_{AB}: \rho \mapsto U\rho U^\dagger$ is at least $1 + 2R(J(\mathcal{U}))$.

If in addition $\mathcal{E}_j = \mathcal{V}_j \otimes \mathcal{W}_j$ for every $j \in [m]$ we call this a **local space-like cut**.

Theorem [PS'22]: The minimal 1-norm of a space-like cut of $\mathcal{U}_{AB}: \rho \mapsto U\rho U^\dagger$ is at least $1 + 2R(J(\mathcal{U}))$.

$$\begin{array}{c} A \\ B \end{array} \begin{array}{|c|} \hline \mathcal{U} \\ \hline \end{array} = \sum_{j=1}^m a_j \begin{array}{c} |0\rangle\langle 0| \text{---} A \\ \text{---} B \\ |0\rangle\langle 0| \text{---} B \end{array} \begin{array}{|c|} \hline \mathcal{V}_j \\ \hline \mathcal{W}_j \\ \hline \end{array} \begin{array}{c} \text{---} O_j^A \\ \text{---} O_j^B \end{array}$$

(Local space-like cut)

Theorem: Let $U = \sum_{\alpha} c_{\alpha} V_{\alpha} \otimes W_{\alpha}$ be a decomposition of U into local unitaries. The **double Hadamard test** is a local space-like cut of $\mathcal{U}: \rho \mapsto U\rho U^\dagger$ with two ancilla qubits and 1-norm $\|a\|_1 = 2\|c\|_1^2 - \|c\|_2^2$. Furthermore, whenever this decomposition is an **operator Schmidt decomposition**,

$$\|a\|_1 = 2\|c\|_1^2 - 1 = 1 + 2R(J(\mathcal{U})).$$

We define the **product extent** $\xi(U)$ as the minimum of $2\|c\|_1^2 - \|c\|_2^2$ over all $U = \sum_{\alpha} c_{\alpha} V_{\alpha} \otimes W_{\alpha}$.

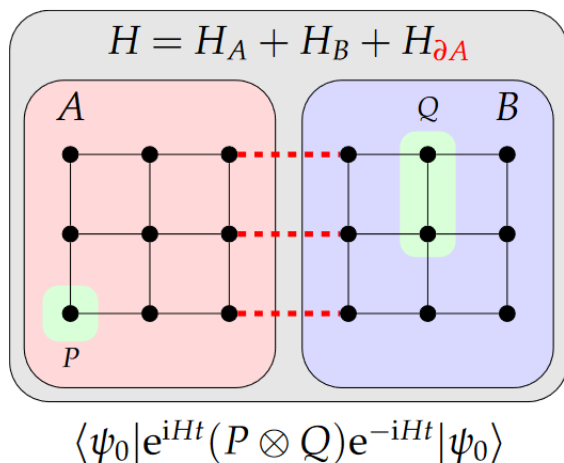
The product extent

We define the **product extent** $\xi(U)$ as the minimum of $2\|c\|_1^2 - \|c\|_2^2$ over all $U = \sum_{\alpha} c_{\alpha} V_{\alpha} \otimes W_{\alpha}$.

The product extent is an entanglement measure satisfying:

1. Faithfulness: $\xi(U) = 1$ iff U is a product of local unitaries.
2. Local unitary invariance: $\xi((V_A \otimes V_B)U(W_A \otimes W_B)) = \xi(U)$.
3. Submultiplicativity: $\xi(UV) \leq \xi(U)\xi(V)$.

Clustered Hamiltonian simulation



Consider the 1st order Trotter formula:

$$U = (\prod_{j \in A} e^{-iH_j t/r} \prod_{k \in B} e^{-iH_k t/r} \prod_{\ell \in \partial A} e^{-iH_\ell t/r})^r.$$

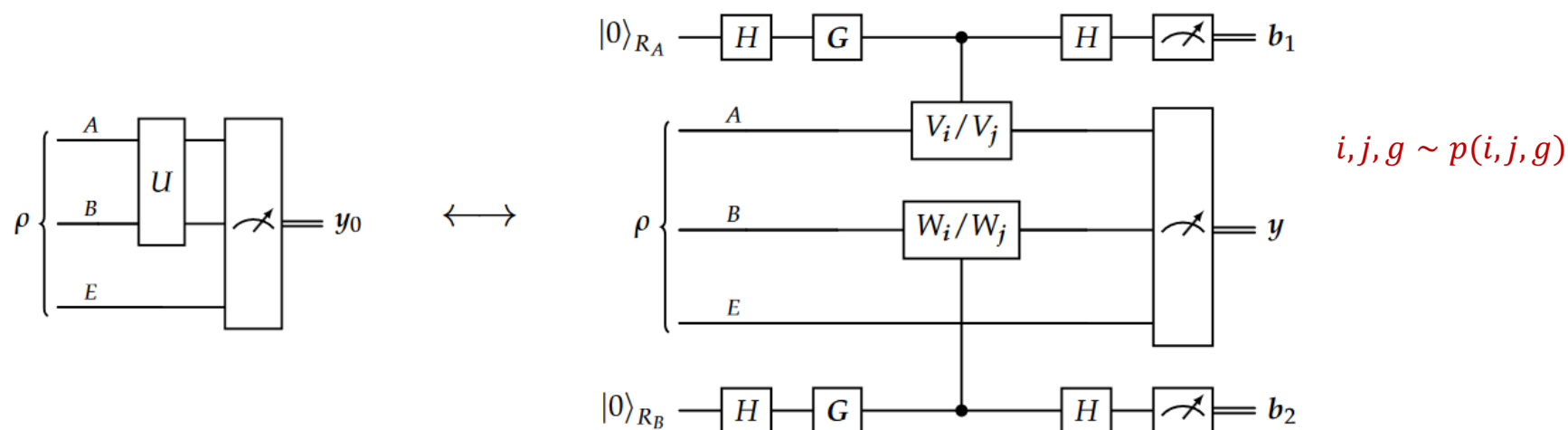
$$\xi(U) \leq (\prod_{j \in A} \xi(e^{-iH_j t/r}) \prod_{k \in B} \xi(e^{-iH_k t/r}) \prod_{\ell \in \partial A} \xi(e^{-iH_\ell t/r}))^r$$

$$= \left(\prod_{\ell \in \partial A} [1 + 2R(J(e^{-iH_\ell t/r}))] \right)^r = 1$$

$$= \left(\prod_{\ell \in \partial A} [1 + |4 \cos(\|H_\ell\|t/r) \sin(\|H_\ell\|t/r)|] \right)^r \leq \exp \left\{ 4 \left(\sum_{\ell \in \partial A} \|H_\ell\| \right) t \right\}$$

The double Hadamard test

Let $U = \sum_{\alpha} c_{\alpha} V_{\alpha} \otimes W_{\alpha}$. We can explain the simulation procedure without referring to QPDs.



Lemma: $\xi(U)(-1)^{b_1+b_2+g}y$ is an unbiased estimator of $\mathbb{E} y_0$.

→ If $\text{Var}[y_0] = \sigma^2$ then the variance of this estimator is $\xi(U)^2 \sigma^2$.

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Time-like cuts

$$\frac{n}{\text{---}} = \sum_{j=1}^m a_j \text{---} \boxed{M_j} y \quad \sigma_y \text{---}$$

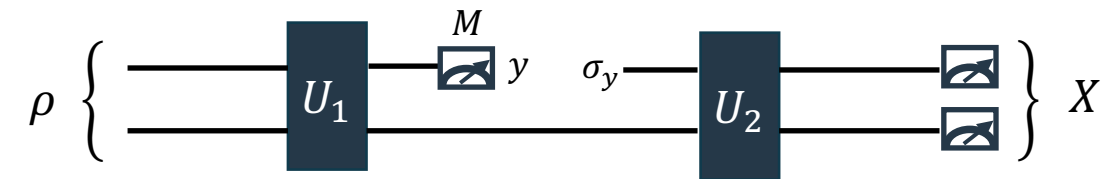
$$\Leftrightarrow id = \sum_{j=1}^m a_j \mathcal{M}_j \quad \leftarrow \text{Measure-and-prepare channels}$$

Theorem [Yua+21, This work]: The optimal time-like cut has 1-norm $2^{n+1} - 1$. Moreover, the measure-and-prepare channels can each be implemented using $O(n^2)$ diagonal gates.

$$id = 2^n \mathcal{M}_0 - (2^n - 1) \mathcal{M}_1$$

How well does this work for a concrete estimation task?

Task: Given fixed U_1, U_2 along with N copies of some unknown state ρ , estimate $\text{Tr}(XU_2U_1\rho U_1^\dagger U_2^\dagger)$ using circuits of the form



If we only care about rank-1 observables, $O(2^n)$ copies suffice.

Theorem: For rank-1 observables, $\Omega(2^n)$ copies are necessary.

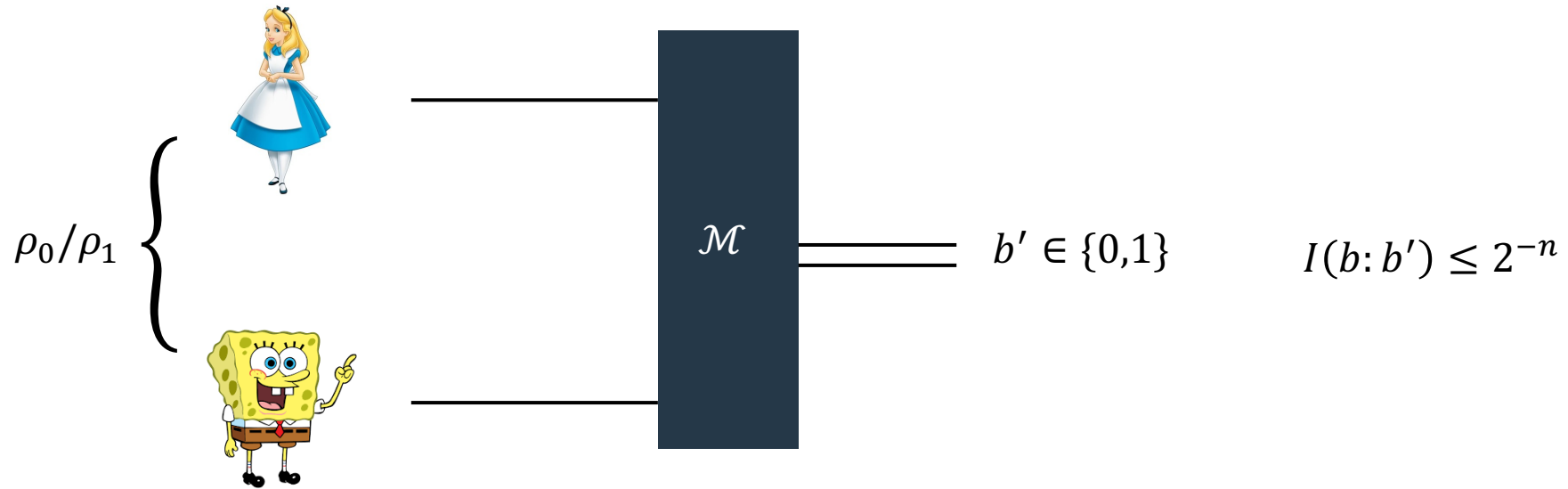
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Lower bounds

Use **quantum data hiding** [DLT'02]: there exists a pair of orthogonal states $(\rho_b: b \in \{0,1\})$ on AB such that b is inaccessible to Alice and Bob if they only make LOCC measurements.



Time-like cut that works too well $\Rightarrow I(b': b) > 2^{-n} \Rightarrow$ contradiction!

Open questions

- Can we find an example of a U such that $\xi(U) \neq 1 + 2R(J(U))$?
- What are some lower bounds in the space-like case?
- Can we extend the clustered Hamiltonian simulation algorithm to compute correlation functions in thermal states?
- Is there a classically-hard family of n -qubit circuits that can be simulated in time $\text{poly}(n)$ using just $\text{polylog}(n)$ -qubit circuits?

- Can a quantum computer outperform classical computers at some useful task?
→ Central motivating question
 - Now (NISQ VQEs), Far future (Fault-tolerance) <https://arxiv.org/pdf/2203.17181.pdf> has a graph cartoon
 - Interesting to look at “soon” regime (hundreds of logical qubits). Quantum simulation?
 - Highlight Quera paper as example of status quo
- Maintaining long-range entanglement is a significant barrier to large-scale quantum computation. How do we mitigate this?
 - Error mitigation: try our best with noise, mitigate the errors
 - Circuit cutting idea: we don’t do the entangling operation at all.
 - Can we apply circuit cutting to practical problems?
- Philosophical question: quantum Shannon theory has been (somewhat) successful in establishing a tradeoff between classical randomness and entanglement, for communication. Is there a similar tradeoff between these resources for computation? Circuit cutting is one answer.